

Non-Parallel Electric and Magnetic Fields in a Gravitational Background, Stationary Gravitational Waves and Gravitons

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Abstract

The existence of an electromagnetic field with parallel electric and magnetic components is readdressed in the presence of a gravitational field. A non-parallel solution is shown to exist. Next, we analyse the possibility of finding stationary gravitational waves in de nature. Finally, We construct a $D=4$ effective quantum gravity model. Tree-level unitarity is verified.

1 Electric and Magnetic Field in a Gravitational Background

Based on a series of papers by Brownstein [1] and Salingaros [2], we readdress here the possibility of the existence of an electromagnetic field whose electric and magnetic components are parallel in the presence of a gravitational field. The coupling between the electromagnetic sector and the gravitational background is accomplished by means of the action.

$$\mathcal{S} = \int \sqrt{-\tilde{g}} \left(-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \right) d^4x, \quad (1.1)$$

where

$$\tilde{g} = \det(g_{\mu\nu}),$$

and

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu.$$

From the above action, the following field-equations follow:

$$\mathcal{D}_\mu F^{\mu\nu} = \partial_\mu F^{\mu\nu} + \Gamma_{\beta\lambda}^\beta F^{\lambda\nu} + \Gamma_{\mu\lambda}^\nu F^{\mu\lambda} = J^\nu \quad (1.2)$$

$$\mathcal{D}_\mu F_{\nu\beta} + \mathcal{D}_\nu F_{\beta\mu} + \mathcal{D}_\beta F_{\mu\nu} = 0. \quad (1.3)$$

Choosing the background to be described by the F.R.W metric,

$$dS^2 = dt^2 - a^2(t) \left[\frac{dr^2}{1 - Ar^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right], \quad (1.4)$$

the Maxwell equations in the absence of electromagnetic sources read as below:

$$\begin{aligned} \vec{\nabla} \cdot \vec{E} &= g \vec{\nabla} f \cdot \vec{E} & (1.5) \\ \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} \\ \vec{\nabla} \times \vec{B} &= \frac{\partial \vec{E}}{\partial t} - g \frac{\partial f}{\partial t} \vec{E} + g \vec{\nabla} f \times \vec{B} - \Gamma_{\mu\beta}^i F^{\mu\beta}, \end{aligned}$$

where

$$g = \frac{\sqrt{1 - Ar^2}}{a^3 r^2 \sin \theta}, \quad f = \frac{a^3 r^2 \sin \theta}{\sqrt{1 - Ar^2}}, \quad (1.6)$$

and $A = +1, 0, -1$.

The Wave-equations for \vec{E} and \vec{B} are found to be given by:

$$\nabla^2 \vec{E} - \frac{\partial^2 \vec{E}}{\partial t^2} = \vec{\nabla} (g \vec{\nabla} f \cdot \vec{E}) - \frac{\partial g}{\partial t} \frac{\partial f}{\partial t} \vec{E} - g \frac{\partial^2 f}{\partial t^2} \vec{E} - g \frac{\partial f}{\partial t} \frac{\partial \vec{E}}{\partial t} + \quad (1.7)$$

$$+ \frac{\partial g}{\partial t} \vec{\nabla} f \times \vec{B} + g \frac{\partial}{\partial t} \vec{\nabla} f \times \vec{B} + g \vec{\nabla} f \times \frac{\partial \vec{B}}{\partial t} - \frac{\partial}{\partial t} (\Gamma_{\mu\beta}^i F^{\mu\beta});$$

$$\nabla^2 \vec{B} - \frac{\partial^2 \vec{B}}{\partial t^2} = \vec{\nabla} \times (g \frac{\partial f}{\partial t} \vec{E} - g \vec{\nabla} f \times \vec{B} + \Gamma_{\mu\beta}^i F^{\mu\beta}) \quad (1.8)$$

Now, by virtue of the presence of the gravitational background, we have explicitly built up a solution for \vec{E} and \vec{B} that evade the claim by Brownstein [1] and Salingaros [2]. These authors state that it is always possible to find out parallel solutions for \vec{E} and \vec{B} in Plasma Physics or in an Astrophysic Plasma. However, contrary to their result, we have found non-parallel solutions due to the non-flat background of gravity:

$$\begin{aligned} \vec{E} = & \hat{i}(\sin \theta G_{(r,t,\theta)} - \cos \theta F_{(r,t)}) ka \cos(kz) \cos(\omega t) + \quad (1.9) \\ & + \hat{j}(\cos \theta F_{(r,t)} - \sin \theta G_{(r,t,\theta)}) ka \sin(kz) \cos(\omega t) + \\ & + \hat{k}[(\sin \theta \cos \varphi F_{(r,t)} + \cos \theta \cos \varphi G_{(r,t,\theta)}) ka \cos(kz) \cos(\omega t) + \\ & - (\sin \theta \sin \varphi F_{(r,t)} + \cos \theta \sin \varphi G_{(r,t,\theta)}) ka \sin(kz) \cos(\omega t)] \end{aligned}$$

and

$$\vec{B} = ka[\hat{i} \sin(kz) + \hat{j} \cos(kz)] \cos(\omega t) \quad (1.10)$$

where the functions $G_{(r,t,\theta)} = \frac{a \cot \theta}{3 \dot{a} r}$ and

$$F_{(r,t)} = \frac{2a}{3\dot{a}r} + \frac{Aar}{3\dot{a}(1 - Ar^2)} \quad (1.11)$$

are the metric contribution.

2 Stationary Gravitational Waves and Gravitons

Now, we analyse the possibility of finding stationary gravitational waves. From a phenomenological viewpoint, a distribution of black holes could play the role of knots for the non-propagating gravitational waves. We postulate the equation that may lead to this sort of waves:

$$R_{\mu\nu} = \kappa\Lambda h_{\mu\nu}, \quad (2.12)$$

$$g_{\mu\nu}(x) = \eta_{\mu\nu} + \kappa h_{\mu\nu}, \quad (2.13)$$

where Λ is the cosmological constant. These equations yield:

$$\partial_\beta\partial_\nu h_\mu^\beta + \partial_\beta\partial_\mu h_\nu^\beta - \square h_{\mu\nu} - \partial_\mu\partial_\nu h_\beta^\beta = \Lambda h_{\mu\nu}. \quad (2.14)$$

Now, solution of the form

$$h_{\mu\nu} = C_{\mu\nu}(z)f(t), \quad (2.15)$$

$$h_{\mu\nu} = \begin{pmatrix} A_{00} & 0 & 0 & 0 \\ 0 & A_{11} & A_{12} & 0 \\ 0 & A_{12} & -A_{11} & 0 \\ 0 & 0 & 0 & A_{00} \end{pmatrix} e^{i\tilde{k}z} \cos \omega t, \quad (2.16)$$

can be found, where A_{00} , A_{11} and A_{12} are free parameters, whereas $\tilde{k} = \sqrt{\Lambda - \omega^2}$ is the wave number. Having in mind that Λ is a small number, the frequency ω must be extremely small. This forces on us to search for a mechanism to detect such low-frequency stationary waves.

The equations of motion stemming from 2.12 may be derived from the density Lagrangian

$$\begin{aligned} \mathcal{L}_H = & \frac{1}{2} H^{\mu\nu} \square H_{\mu\nu} - \frac{1}{4} H \square H - \frac{1}{2} H^{\mu\nu} \partial_\mu \partial_\alpha H_\nu^\alpha - \frac{1}{2} H^{\mu\nu} \partial_\nu \partial_\alpha H_\mu^\alpha + \\ & - \frac{1}{2} \Lambda H^{\mu\nu} H_{\mu\nu} + \frac{1}{4} \Lambda H^2, \end{aligned} \quad (2.17)$$

where

$$H_\nu^\alpha = h_\nu^\alpha - \frac{1}{2} \delta_\nu^\alpha h$$

and the bilinear form correspondent operator of Lagrangian 2.17 is given by

$$\begin{aligned} \Theta_{\mu\nu,\kappa\lambda} = & (\square - \Lambda)P^{(2)} - \Lambda P_m^{(1)} + \frac{5}{2}(\square - \Lambda)P_s^{(0)} - \frac{(\Lambda + 3\square)}{2}P_w^{(0)} + \quad (2.18) \\ & + \frac{\sqrt{3}}{2}(\Lambda - \square)P_{sw}^{(0)} + \frac{\sqrt{3}}{2}(\Lambda - \square)P_{ws}^{(0)}, \end{aligned}$$

and $P^{(i)}$, $i=0,1,2$, are spin-projetor operators in the space of rank-2 symmetric tensors. The graviton propagator read off from this Lagrangian is given by:

$$\langle T(h_{\mu\nu}(x); h_{\kappa\lambda}(y)) \rangle = i\Theta_{\mu\nu,\kappa\lambda}^{-1} \delta^4(x - y) \quad (2.19)$$

where

$$\Theta^{-1} = \left[X P^{(2)} + Y P_m^{(1)} + Z P_s^{(0)} + W P_w^{(0)} + R P_{sw}^{(0)} + S P_{ws}^{(0)} \right]_{\mu\nu,\kappa\lambda} \quad (2.20)$$

with

$$\begin{aligned} X &= -\frac{1}{\Lambda - \square}; & Y &= -\frac{1}{\Lambda}; & Z &= -\frac{\Lambda + 3\square}{\Lambda^2 + 8\Lambda\square - 9\square^2}; \quad (2.21) \\ W &= -\frac{5}{\Lambda - 9\square}; & R &= -\frac{\sqrt{3}}{\Lambda + 9\square}; & S &= -\frac{\sqrt{3}}{\Lambda + 9\square}; \end{aligned}$$

From this propagator, we can set a current-current amplitude and discuss tree-level unitarity [3]. Three massive excitations were found: They are a spin-2 quantum with mass equal to $k^2 = \Lambda$ and two massive spin-0 quanta with masses equal to $k^2 = \Lambda$ and $k^2 = -\frac{1}{9}\Lambda$. The spin-2 is a physical one: the imaginary part of the residue of the amplitude at the pole $k^2 = \Lambda$ is positive, so that it does not lead to a ghost. It remains to be shown that the tachyonic pole, $k^2 = -\frac{1}{9}\Lambda$, is non-dynamical or decouples through some constraint on the sources.

We conclude, then, that in a gravitational background it is always possible to find non-parallel electric and magnetic fields. It is the gravitational field that breaks the parallel configuration of \vec{E} and \vec{B} argued by [1, 2]. Also, a stationary gravitational wave equation was postulated and a particular solution was found. We argue that such a kind of solution is possible to be found in Black Hole distributions. Finally we set up an effective quantum gravity model where the necessary condition for the tree-level unitarity for the spin-2 sector is respected. The model is infrared finite though non-renormalizable in the ultraviolet limit.

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